

# Axially symmetric lenses

$$\theta \equiv |\theta|$$

see Schneider eqn 40  $\rightarrow K(\theta) = K(\theta)$  (49)

Generalise (10), (21), (22) for axially symmetric lens.  
 Use Birkhoff's theorem: mass enclosed by sphere of radius  $r$  acts gravitationally like a point at the centre, of mass = enclosed mass (for spherically symmetric mass dist<sup>n</sup> only).

Also, lens bend angle does not care exactly how mass is distributed along line of sight if all roughly at  $D_s$ .

$\rightarrow$  in eqn (1) this can be mass enclosed at radius  $r$  follow this through derivation of eqn (10)  $\rightarrow$

(10) becomes 
$$\alpha = \frac{D_{ds}}{D_s D_d} \frac{4G}{c^2} \frac{\theta}{|\theta|^2} M_{\text{enc}} \quad (50)$$

plus use (5): 
$$\alpha = \frac{D_{ds}}{D_s D_d} \frac{4G}{c^2} \frac{1}{|\theta|} M_{\text{enc}} \quad (51)$$

Using notation in (36)  $|\theta'| < |\theta|$

$\rightarrow M_{\text{enc}}(\theta) \equiv M_{\text{enc}}(\theta) = \int_{|\theta'|=0}^{\theta} dm = \int \Sigma(D_u \theta) D_u^2 d\theta'$

$$= \sum_{\text{cst}} D_d^2 \int_{\theta'=0}^{\theta} K(\theta') 2\pi \theta' d\theta' \quad (52)$$

$d\theta'$  becomes thin  
 Since now lets integrate in azimuthal (e.g. Schneider eqn 34) (53)

$\rightarrow \alpha = \frac{2}{|\theta|} \int_{\theta'=0}^{\theta} K(\theta') \theta' d\theta'$

use  $\int_{\text{cst}} d\theta'$